

Eigenfunctions on a Riemannian manifold and representations of a vertex operator algebra

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Field Theory"

Outline

The vertex operator algebra approach to two-dimensional conformal field theory

Nonlinear sigma models

Classical nonlinear sigma models: harmonic maps
Quantization

Heisenberg algebras and smooth functions

A homomorphism from the symmetric algebra $S(T_p M)$ to the algebra of linear operators on $C^\infty(B(p, r))$

Modules for the Heisenberg algebra $\widehat{T_p M}$ generated by germs of smooth functions at p

$L(0)$ -semisimple lower-bounded generalized modules generated by eigenfunctions

Sheaves of vertex operator algebras and weak modules
Eigenfunctions and $L(0)$ -semisimple lower-bounded generalized modules for a vertex operator algebra

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Vertex operator algebras

Belavin-Polyakov-Zamolodchikov, Borchers,
Frenkel-Lepowsky-Meurman.

A **vertex operator algebra** consists the following data:

- ▶ A \mathbb{Z} -graded vector space $V = \coprod_{n \in \mathbb{Z}} V_{(n)}$.
- ▶ A **vertex operator map**

$$\begin{aligned} Y_V : V \otimes V &\rightarrow V[[z, z^{-1}]], \\ u \otimes v &\mapsto Y(u, z)v. \end{aligned}$$

- ▶ A **vacuum** $\mathbf{1} \in V_{(0)}$.
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These data satisfy the following axioms:

- ▶ **Grading-restriction property:** $\dim V_{(n)} < \infty$ for $n \in \mathbb{Z}$ and $V_{(n)} = 0$ when n is sufficiently negative.
- ▶ **Lower-truncation property:** For $u, v \in V$, $Y(u, z)v$ contains only finitely many negative power terms.
- ▶ **Axioms for the vacuum:** For $u \in V$, $Y(\mathbf{1}, z)u = u$ and $\lim_{z \rightarrow 0} Y(u, z)\mathbf{1} = u$.
- ▶ **Axioms for the conformal element:** Let $L(n) : V \rightarrow V$ be defined by $Y(\omega, z) = \sum_{n \in \mathbb{Z}} L(n)z^{-n-2}$, then $[L(m), L(n)] = (m - n)L(m + n) + \frac{c}{12}(m^3 - m)\delta_{m+n, 0}$, $\frac{d}{dz} Y(u, z) = Y(L(-1)u, z)$ (**$L(-1)$ -derivative property**) for $u \in V$ and $L(0)u = nu$ for $u \in V_{(n)}$ (**$L(0)$ -grading property**).

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Vertex operator algebras

- ▶ **Duality property:** For $u_1, u_2, v \in V, v' \in V' = \coprod_{n \in \mathbb{Z}} V_{(n)}^*$, the series

$$\begin{aligned} & \langle v', Y(u_1, z_1) Y(u_2, z_2) v \rangle \\ & \langle v', Y(u_2, z_2) Y(u_1, z_1) v \rangle \\ & \langle v', Y(Y(u_1, z_1 - z_2) u_2, z_2) v \rangle \end{aligned}$$

are absolutely convergent in the regions $|z_1| > |z_2| > 0$, $|z_2| > |z_1| > 0$ and $|z_2| > |z_1 - z_2| > 0$, respectively, to a common rational function in z_1 and z_2 with the only possible poles at $z_1, z_2 = 0$ and $z_1 = z_2$.

Modules

- ▶ Let V be a vertex operator algebra. A **V -module** is a \mathbb{C} -graded vector space $W = \coprod_{n \in \mathbb{C}} W_{(n)}$ equipped with a vertex operator map $Y_W : V \otimes W \rightarrow W[[z, z^{-1}]]$ satisfying all those axioms for V which still make sense.
- ▶ A **lower-bounded generalized V -module** is a \mathbb{C} -graded vector space $W = \coprod_{n \in \mathbb{C}} W_{(n)}$ equipped with a vertex operator map $Y_W : V \otimes W \rightarrow W[[z, z^{-1}]]$ satisfying all the axioms for a V -module except that $W_{(n)}$ are generalized eigenspaces instead of eigenspaces of $L(0)$ and are not necessarily finite dimensional. A **weak V -module** is a vector space W equipped with a vertex operator map $Y_W : V \otimes W \rightarrow W[[z, z^{-1}]]$ satisfying all those axioms for V -modules which still make sense.

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Intertwining operators

Let W_1 , W_2 and W_3 be V -modules. An **intertwining operator** of type $\begin{pmatrix} W_3 \\ W_1 W_2 \end{pmatrix}$ is a linear map $\mathcal{Y} : W_1 \otimes W_2 \rightarrow W_3\{z\}$, where $W_3\{z\}$ is the space of all series in complex powers of z with coefficients in W_3 , satisfying all those axioms for V which still make sense, that is, a lower-truncation property, an $L(-1)$ -derivative property and a duality property.

Examples

- ▶ Free bosons: Representations of infinite-dimensional Heisenberg algebras.
- ▶ Free bosons on tori: Vertex operator algebras, modules and classical vertex operators (intertwining operators) associated to lattices.
- ▶ Wess-Zumino-Novikov-Witten models: Representations of affine Lie algebras.
- ▶ Minimal models: Representations of the Virasoro algebra.
- ▶ Fermion theories: Representations of infinite-dimensional Clifford algebras, affine Lie superalgebras and superconformal algebras.
- ▶ Orbifolds, cosets and \mathcal{W} -algebras, including in particular the moonshine module.

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Rational conformal field theories

Theorem (H. 2002–2005)

Let V be a simple vertex operator algebra satisfying the following conditions:

- 1. $V_{(n)} = 0$ for $n < 0$, $V_{(0)} = \mathbb{C}\mathbf{1}$ and V' is isomorphic to V as a V -module.*
- 2. Every lower-bounded generalized V -module is a direct sum of irreducible V -modules.*
- 3. V is C_2 -cofinite, that is, $\dim V/C_2(V) < \infty$ where $C_2(V)$ is the subspace of V spanned by elements of the form $\text{Res}_z z^{-2} Y(u, z)v$ for $u, v \in V$.*

Then we have the following conclusion:

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Then we have the following conclusion:

1. *The Moore-Seiberg polynomial equations hold. In particular, the Verlinde conjecture holds.*
2. *The category of V -modules has a natural structure of a modular tensor category. In particular, we have a modular functor (in all genera) and a 3-dimensional topological field theory.*
3. *All chiral and full correlation functions in genus-zero and genus-one can be constructed from intertwining operators (the part on full correlation functions being done jointly with Kong).*
4. *There exist locally convex topological completions of the spaces involved such that Segal's axioms in genus-zero and genus-one are satisfied.*

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- ▶ Many results for rational theories can be generalized to non-rational theories, including unitary non-rational theories and logarithmic theories. But there are still conjectures and open problems to be solved. The construction and study of non-rational theories is important for the study of the moduli space of conformal field theories. Nonlinear sigma models with Calabi-Yau manifolds as targets are important examples to be constructed and studied.

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Harmonic maps from a Riemann surface to a Riemannian manifold

- ▶ Σ : 2-dimensional Riemannian manifold.
 M : n -dimensional Riemannian manifold.
 φ : smooth map from Σ to M .

- ▶ Action: $\int_{\Sigma} \|d\varphi\| dS$.

Locally: $\|d\varphi\| = \eta^{ij} g_{\mu\nu} \frac{\partial\varphi^{\mu}}{\partial x^i} \frac{\partial\varphi^{\nu}}{\partial x^j}$.

- ▶ Harmonic maps: Critical points of the action above.
- ▶ The action is conformally invariant. Only the Riemann surface structure on Σ is needed.

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The vertex operator algebra approach to two-dimensional conformal field theory

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Classical nonlinear sigma models: harmonic maps

Quantization

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Free bosons: Quantization of linear sigma models

- ▶ Linear sigma models: M flat.
- ▶ $M = \mathbb{R}^n$ and $p \in \mathbb{R}^n$.
 $T_p M = T_p \mathbb{R}^n$: the tangent space of $M = \mathbb{R}^n$ at p .
- ▶ Heisenberg algebra: $\widehat{T_p M} = \widehat{T_p M}_- \oplus \widehat{T_p M}_0 \oplus \widehat{T_p M}_+$, where

$$\widehat{T_p M}_- = T_p M \otimes t^{-1} \mathbb{C}[t^{-1}],$$

$$\widehat{T_p M}_0 = T_p M \otimes \mathbb{C}t^0 \oplus \mathbb{C}I \simeq T_p M \oplus \mathbb{C}I,$$

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- ▶ Eigenfunctions $|\mathbf{k}\rangle = e^{ik_\mu x^\mu}$ of the Laplacian with eigenvalues $-k_\mu k^\mu$.

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Vertex operator algebra and modules associated to free bosons

- ▶ Let $\widehat{T}_\rho M_+$ act on $|\mathbf{k}\rangle$ as 0, \mathbf{l} on $|\mathbf{k}\rangle$ as 1 and elements of $T_\rho M$ on $|\mathbf{k}\rangle$ as i times the vector fields obtained by parallel transporting the elements of $T_\rho M$. Then $\mathbb{C}|\mathbf{k}\rangle$ becomes a module for $\widehat{T}_\rho M_0 \oplus \widehat{T}_\rho M_+$.
- ▶ Induced module: $W_{|\mathbf{k}\rangle} = U(\widehat{T}_\rho M) \otimes_{U(\widehat{T}_\rho M_0 \oplus \widehat{T}_\rho M_+)} \mathbb{C}|\mathbf{k}\rangle$
(Fock space generated by $|\mathbf{k}\rangle$).
- ▶ $W_{|0\rangle}$ is a vertex operator algebra and $W_{|\mathbf{k}\rangle}$ for any \mathbf{k} is a module for the vertex operator algebra $W_{|0\rangle}$.
- ▶ The corresponding conformal field theory can be constructed by studying the correlation functions among these modules for the vertex operator algebra $W_{|0\rangle}$.

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Sigma models on tori

- ▶ Given a torus, the vertex operator algebra is still $W_{|0\rangle}$. But consider only modules $W_{|\mathbf{k}\rangle}$ for those \mathbf{k} such that $e^{ik_\mu x^\mu}$ are well defined on the torus. Such \mathbf{k} form a lattice.
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First difficulty in generalizing this construction to the curved case

- ▶ M : a Riemannian manifold. $p \in M$. T_pM : the tangent space of \mathbb{R}^n at p .
- ▶ Heisenberg algebra $\widehat{T_pM}$.
- ▶ Problem: How to define an action of the abelian Lie subalgebra T_pM on smooth functions, in particular on eigenfunctions of the Laplacian, at least near p .
- ▶ If we choose a local coordinate patch near p and define the action using derivatives with respect to the coordinates, the results depend on the local coordinate patch.
- ▶ If we define the action using the covariant derivatives, we do not have a representation of the abelian Lie subalgebra T_pM because the commutators of covariant derivatives give the curvature tensor.

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Notations

- ▶ U : an open subset of M .
 $\Gamma_U(S(TM))$: the space of sections on U of the vector bundle $S(TM)$ whose fibers are the symmetric algebras of tangent spaces of M .
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Definition of η_U

- ▶ $\sigma \in \Gamma_U(\mathcal{S}^k(TM))$ and $q \in U$.
 $\psi : N \rightarrow \mathbb{R}^n$: a chart of normal coordinates at q .
- ▶ Identify the symmetric algebra $\mathcal{S}(T_qM)$ with the space of symmetric tensors at q . Then there exist $a^{\mu_1 \dots \mu_k}(q) \in \mathbb{C}$, symmetric with respect to μ_1, \dots, μ_k , such that

$$\sigma(q) = a^{\mu_1 \dots \mu_k}(q) \left. \partial_{\mu_1}^\psi \right|_q \otimes \dots \otimes \left. \partial_{\mu_k}^\psi \right|_q.$$

- ▶ For $f \in C^\infty(U)$, we define

$$(\eta_U(\sigma)f)(q) = a^{\mu_1 \dots \mu_k}(q) \left. \partial_{\mu_1}^\psi \dots \partial_{\mu_k}^\psi f \right|_q.$$

- ▶ Example: Let g^{-1} be the element of $\Gamma_U(\mathcal{S}^2(TM))$ corresponding to the metric on the cotangent bundle. Then $\eta_U(g^{-1}) = \Delta$.

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A homomorphism of associative algebras

- ▶ $p \in M$ and $r \leq \text{inj}(p)$ (injectivity radius at p).
 $B(p, r)$: open ball of radius r centered at p .
- ▶ Given an element of $S(T_p M)$, using parallel transport along the unique geodesics from p to points in $B(p, r)$, we obtain an element of $\Gamma_{B(p, r)} S(TM)$. Apply $\eta_{B(p, r)}$ to this element, we obtain an operator on $C^\infty(B(p, r))$. Thus we obtain a linear map $\xi_{B(p, r)} : S(T_p M) \rightarrow L(C^\infty(B(p, r)))$.

Theorem

$\xi_{B(p, r)}$ is a homomorphism of associative algebras.

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The space $C^\infty(p)$ of germs of smooth functions at p

- ▶ Let $\widehat{T_p M}_+$ act on $C^\infty(p)$ as 0, $\mathbf{1}$ on $C^\infty(p)$ as 1 and elements of $T_p M$ on $C^\infty(p)$ by i times the homomorphism $\xi_{B(p,r)}$ for suitable r .
- ▶ By the theorem above, $C^\infty(p)$ becomes a module for $\widehat{T_p M}_0 \oplus \widehat{T_p M}_+$.

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Induced module for $\widehat{T}_p M$

- ▶ Induced module:

$$W_{C^\infty(p)} = U(\widehat{T}_p M) \otimes_{U(\widehat{T}_p M_0 \oplus \widehat{T}_p M_+)} C^\infty(p).$$

- ▶ The submodule $W_1(p)$ of $W_{C^\infty(p)}$ generated by the germ containing the constant function 1 is a vertex operator algebra.
- ▶ $W_{C^\infty(p)}$ is a weak module for the vertex operator algebra W_1 . $L(0)$ acts on elements of $C^\infty(p)$ as $-\frac{1}{2}\Delta$.
- ▶ The submodules generated by eigenfunctions of the Laplacians are $L(0)$ -semisimple lower-bounded generalized modules for $W_1(p)$.

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A presheaf

- ▶ U : open subset of M .
 σ : parallel section on U of $U(\widehat{TM})$ (the vector bundle whose fibers are universal enveloping algebras of Heisenberg algebras).
- ▶ Given such σ and $f \in C^\infty(U)$, we obtain a map from U to the disjoint union of
 $W_{C^\infty(p)} = U(\widehat{T_p M}) \otimes_{U(\widehat{T_p M}_0 \oplus \widehat{T_p M}_+)} C^\infty(p)$ for $p \in U$.
- ▶ $W^0(U)$: space spanned by all such maps from U to the disjoint union of $W_{C^\infty(p)}$ for $p \in U$.
- ▶ $U \mapsto W^0(U)$ gives a presheaf \mathcal{W}^0 on M .

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- ▶ U : open subset of M .
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A sheaf of vertex operator algebra and a sheaf of weak modules

- ▶ Let \mathcal{W} be the sheafification of \mathcal{W}^0 .
- ▶ For an open subset U , the subspace $W_1(U)$ of $W^0(U)$ obtained from parallel sections on U of $U(\widehat{TM})$ and the constant function 1 has a structure of a vertex operator algebra. $U \mapsto W_1(U)$ gives a sheaf \mathcal{W}_1 of vertex operator algebras.

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The space $W(U)$ of sections of \mathcal{W} on U is a weak module for $W_1(U)$ and \mathcal{W} is a sheaf of weak modules for the sheaf \mathcal{W}_1 of vertex operator algebras.

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Outline

The vertex operator algebra approach to two-dimensional conformal field theory

Nonlinear sigma models

Classical nonlinear sigma models: harmonic maps
Quantization

Heisenberg algebras and smooth functions

A homomorphism from the symmetric algebra $S(T_p M)$ to the algebra of linear operators on $C^\infty(B(p, r))$

Modules for the Heisenberg algebra $\widehat{T_p M}$ generated by germs of smooth functions at p

$L(0)$ -semisimple lower-bounded generalized modules generated by eigenfunctions

Sheaves of vertex operator algebras and weak modules
Eigenfunctions and $L(0)$ -semisimple lower-bounded generalized modules for a vertex operator algebra

$L(0)$ -semisimple lower-bounded generalized modules for the vertex operator algebra $W_1(M)$ generated by eigenfunctions

Take $U = M$ and let $W_f(M)$ be the weak submodule of $W(M)$ for $W_1(M)$ generated by an eigenfunction f . Then we have:

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$W_f(M)$ is an $L(0)$ -semisimple lower-bounded generalized module for $W_1(M)$.

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The correspondence $M \mapsto (W_1, \{W_f\}_{f \text{ are eigenfunctions}})$ is a functor from the category of Riemannian manifold to the category of pairs of vertex operator algebras and sequences of $L(0)$ -semisimple lower-bounded generalized V -modules.

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